

# STEERING WITNESSES FOR GAUSSIAN STATES

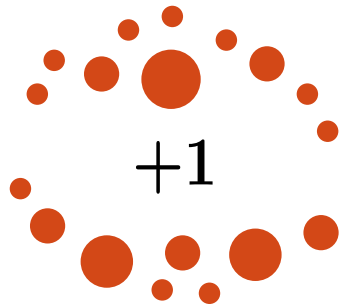
Tatiana Mihaescu<sup>1</sup>, Hermann Kampermann<sup>2</sup>, Aurelian Isar<sup>1</sup>,  
Dagmar Bruss<sup>2</sup>



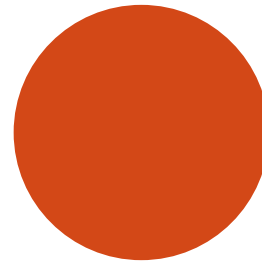
1. National Institute for Physics and Nuclear Engineering, Department of Theoretical Physics, Bucharest, Romania
2. Heinrich-Heine-Universität Dusseldorf, Institut für Theoretische Physik III, Dusseldorf, Germany

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where  $|0\rangle = |z^+\rangle$  and  $|1\rangle = |z^-\rangle$  denote the two possible spin orientations in  $z$ -direction.



Alice:  $z$  direction



Bob's state

## EPR PAPER

Einstein – Podolski – Rosen Paradox

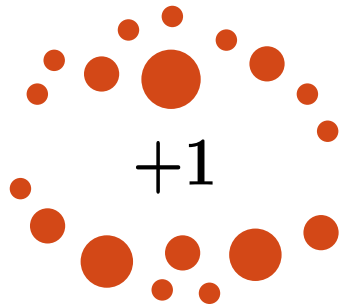
“Can Quantum-Mechanical description of Physical Reality be Considered Complete?”

(1935)

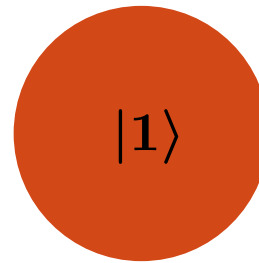


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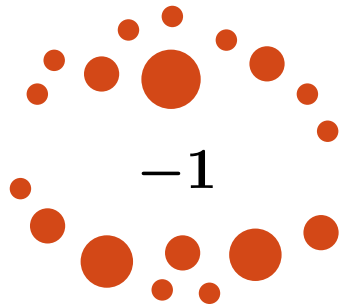
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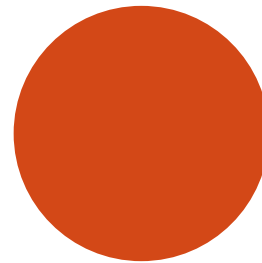


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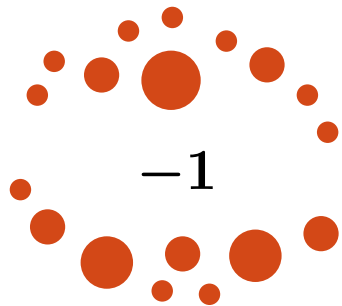
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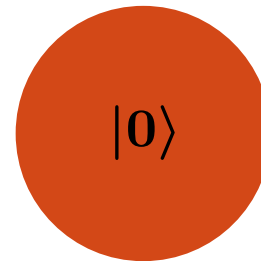


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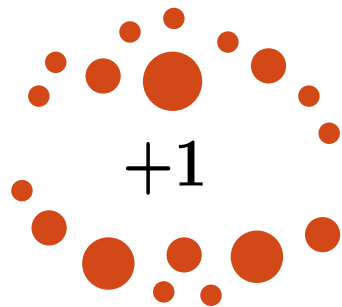
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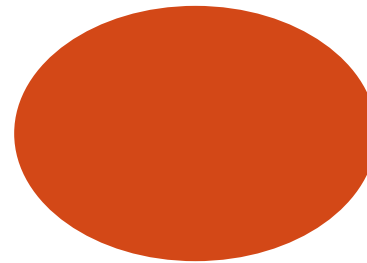


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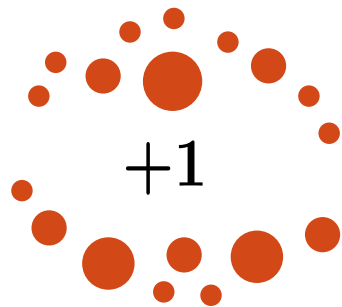
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*Then either (1) the description of reality given by the wave function in quantum mechanics is not complete or (2) these two quantities cannot have simultaneous reality.”*

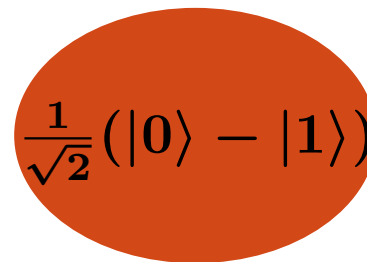


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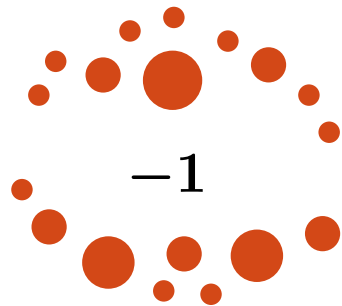
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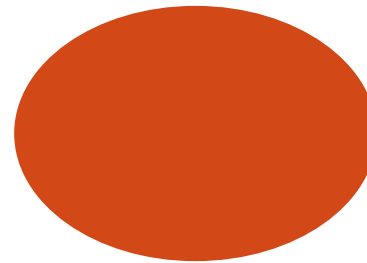


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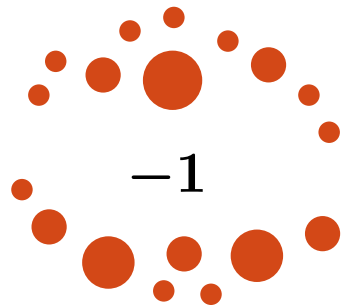
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*“We are thus forced to conclude that the quantum-mechanical description of physical reality given by wave functions is not complete.”*

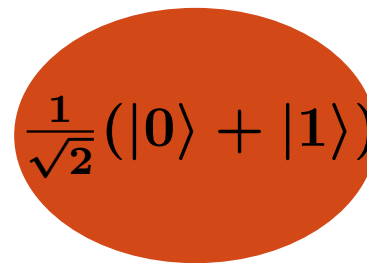


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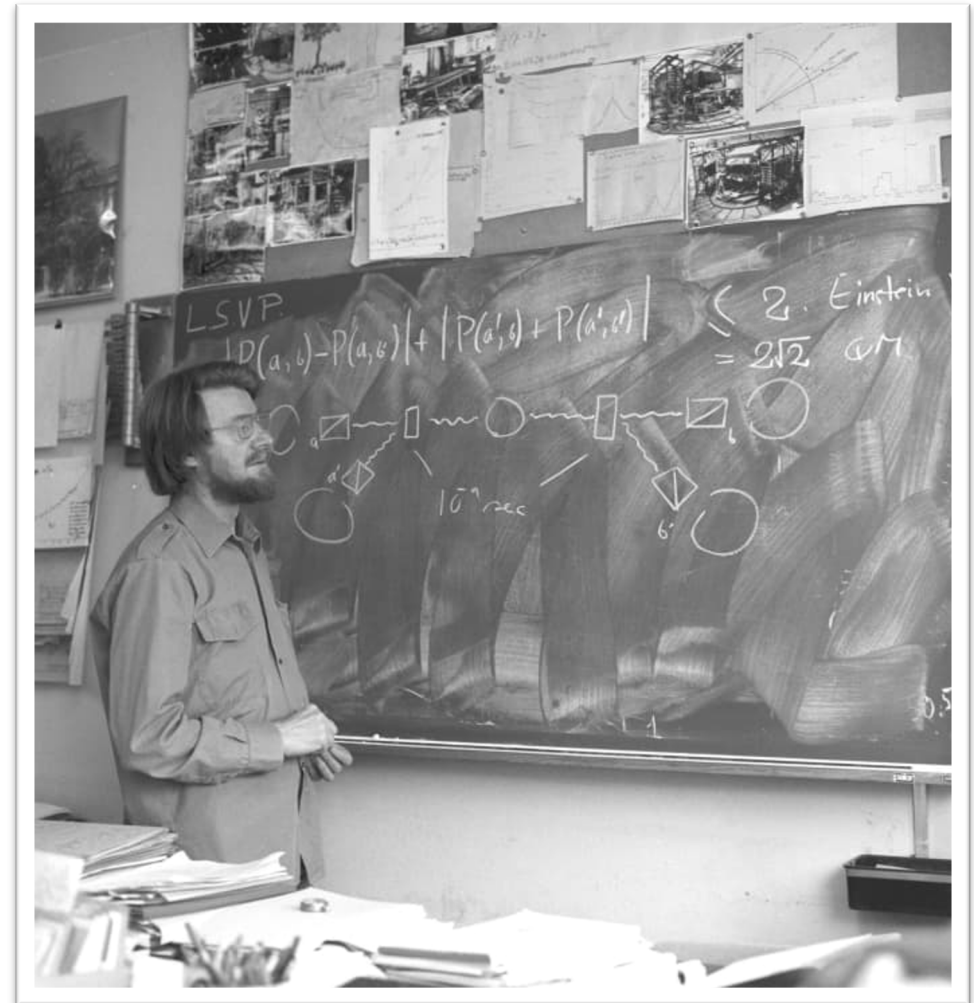






# BELL TEST

*“On the Einstein Podolski Rosen paradox”*  
(1964)

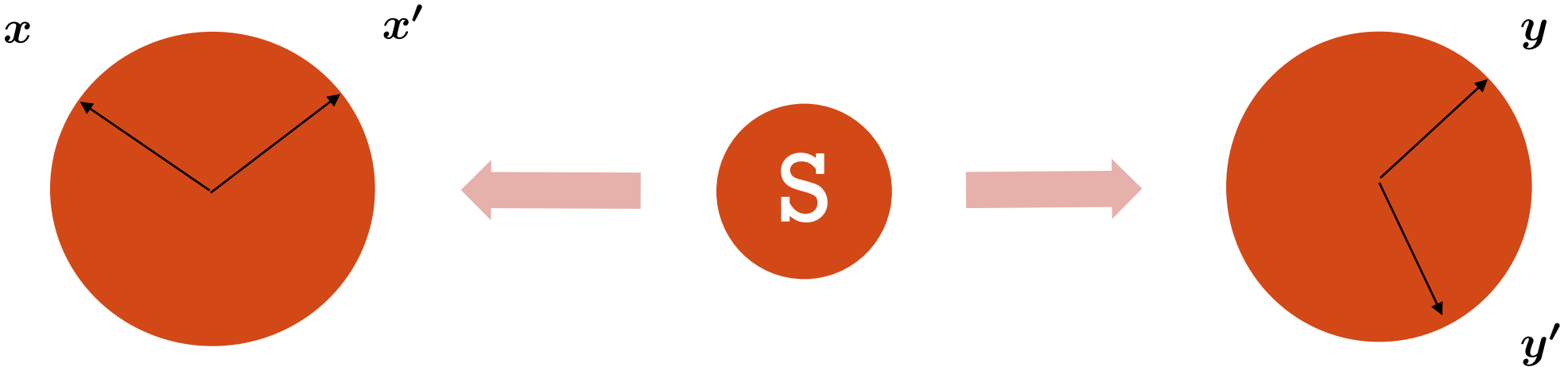


Local Realism?

Local Hidden Variable?



# BELL TEST



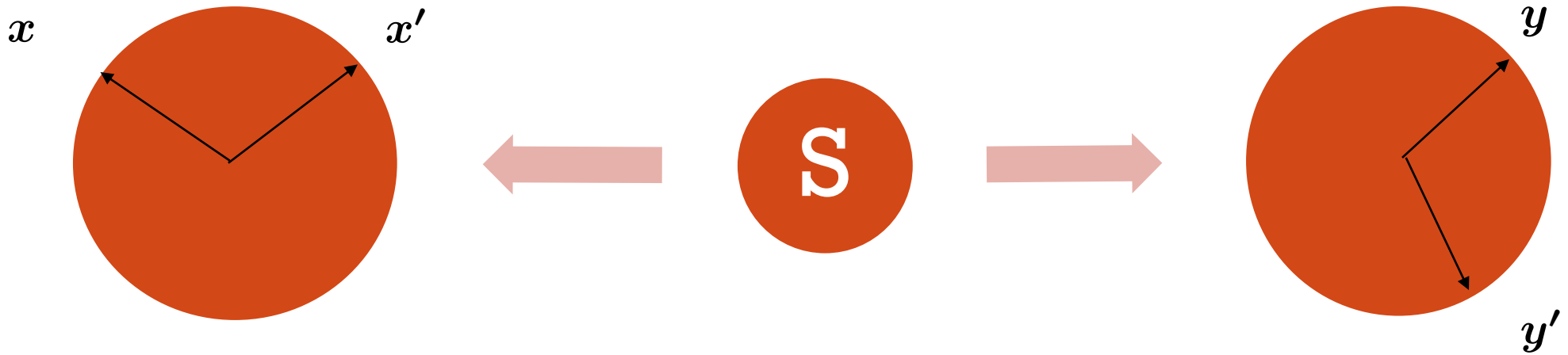
Alice

Measurements:  $A, a \in \{\pm 1\}$   
 $A', a' \in \{\pm 1\}$

Bob

Measurements:  $B, b \in \{\pm 1\}$   
 $B', b' \in \{\pm 1\}$





Outcomes:  $A, a \in \{\pm 1\}$

$A', a' \in \{\pm 1\}$

Outcomes:  $B, b \in \{\pm 1\}$

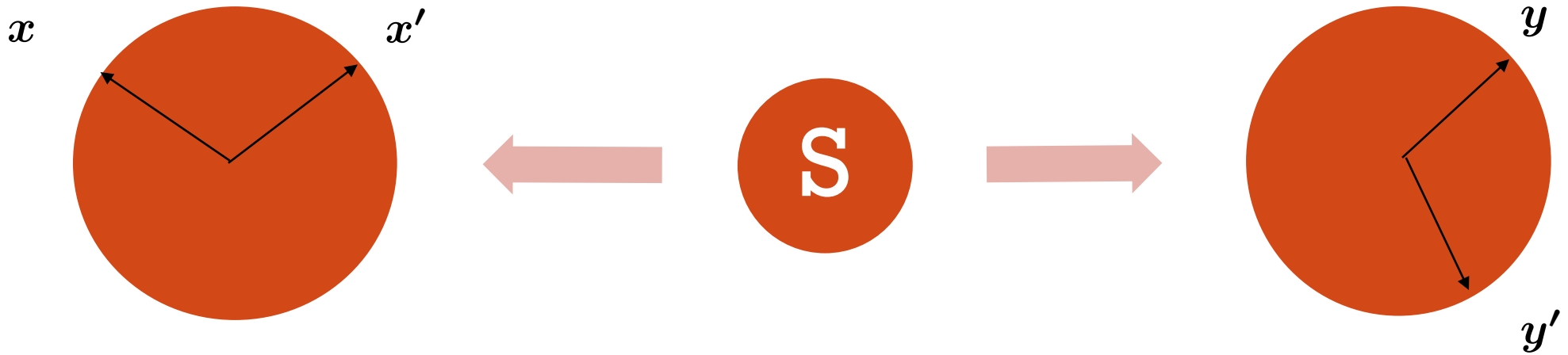
$B', b' \in \{\pm 1\}$

Local Realism:

Quantum mechanics:

<b>a</b>	<b>a'</b>	<b>b</b>	<b>b'</b>
+1	+1	+1	-1
+1	+1	-1	+1
+1	-1	+1	+1
...	...	...	...





Outcomes:  $A, a \in \{\pm 1\}$

$A', a' \in \{\pm 1\}$

Outcomes:  $B, b \in \{\pm 1\}$

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<b>a</b>	<b>a'</b>	<b>b</b>	<b>b'</b>
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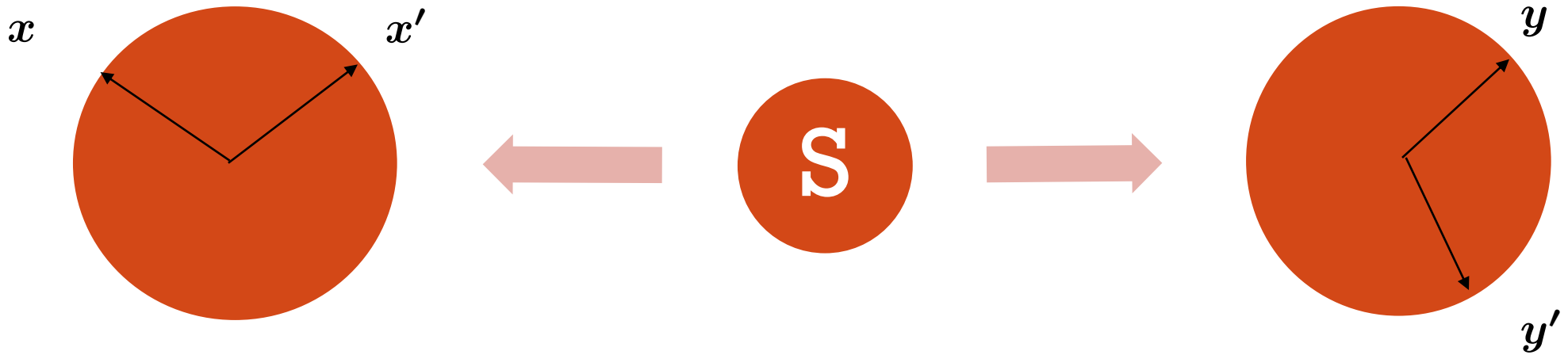
Local Realism:

$$ab + a'b + ab' - a'b' \leq 2$$

Quantum mechanics:







Outcomes:  $A, a \in \{\pm 1\}$   
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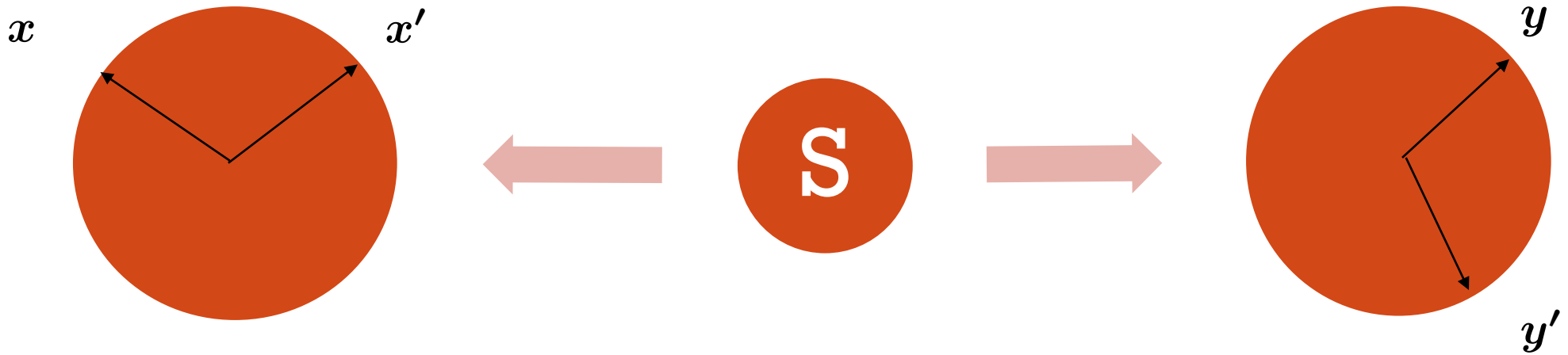
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Quantum mechanics:

$$\langle AB \rangle + \langle A'B \rangle + \langle AB' \rangle - \langle A'B' \rangle \leq 2\sqrt{2}$$





Outcomes:  $A, a \in \{\pm 1\}$   
 $A', a' \in \{\pm 1\}$

Outcomes:  $B, b \in \{\pm 1\}$   
 $B', b' \in \{\pm 1\}$

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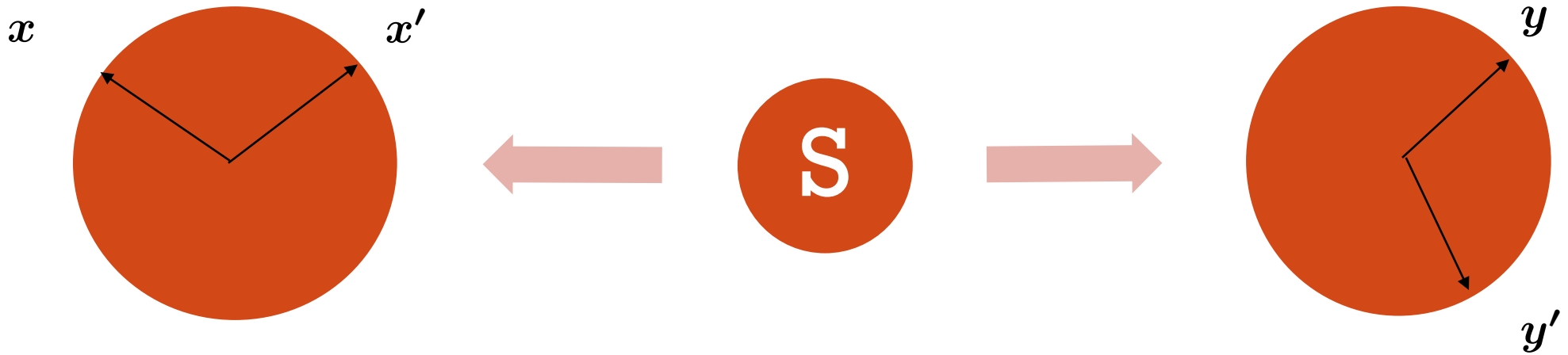
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Einstein was wrong!





# BELL NONLOCALITY

Failure of Local Hidden Variable (LHV) model

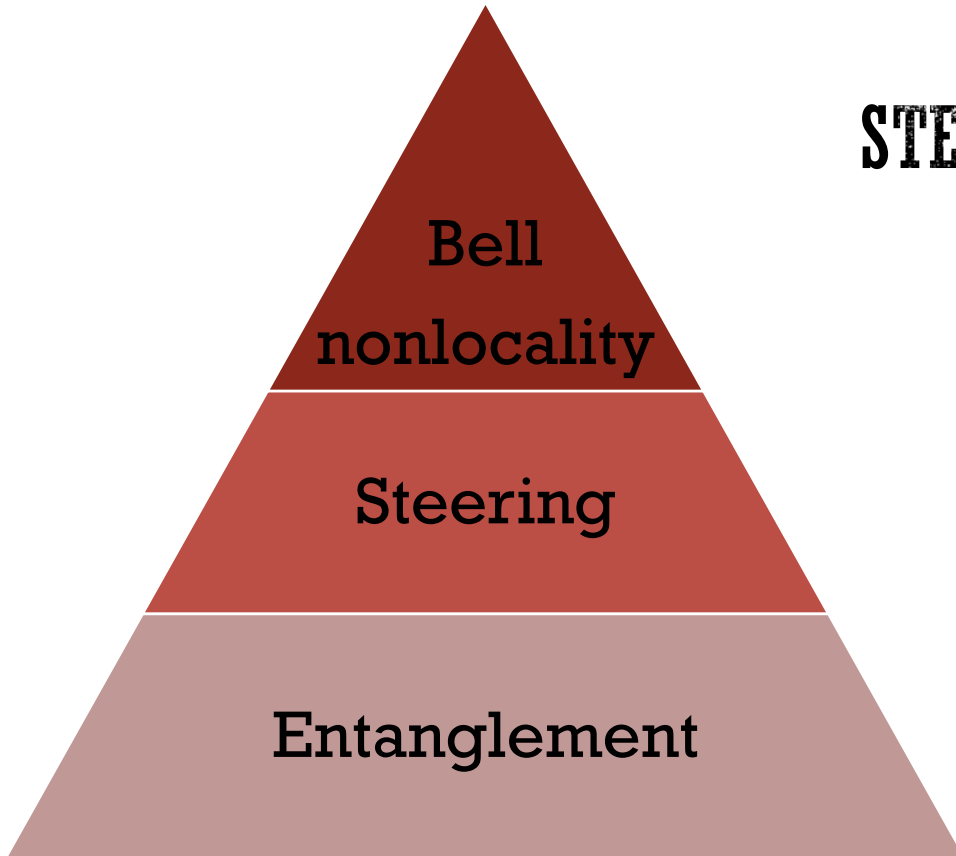
$$p(a, b|x, y) = \int d\lambda p(\lambda) p(a|x, \lambda) p(b|y, \lambda)$$

# STEERING

Failure of Hybrid LHV – LQS model

# ENTANGLEMENT

Failure of Local Quantum State (LQS) model



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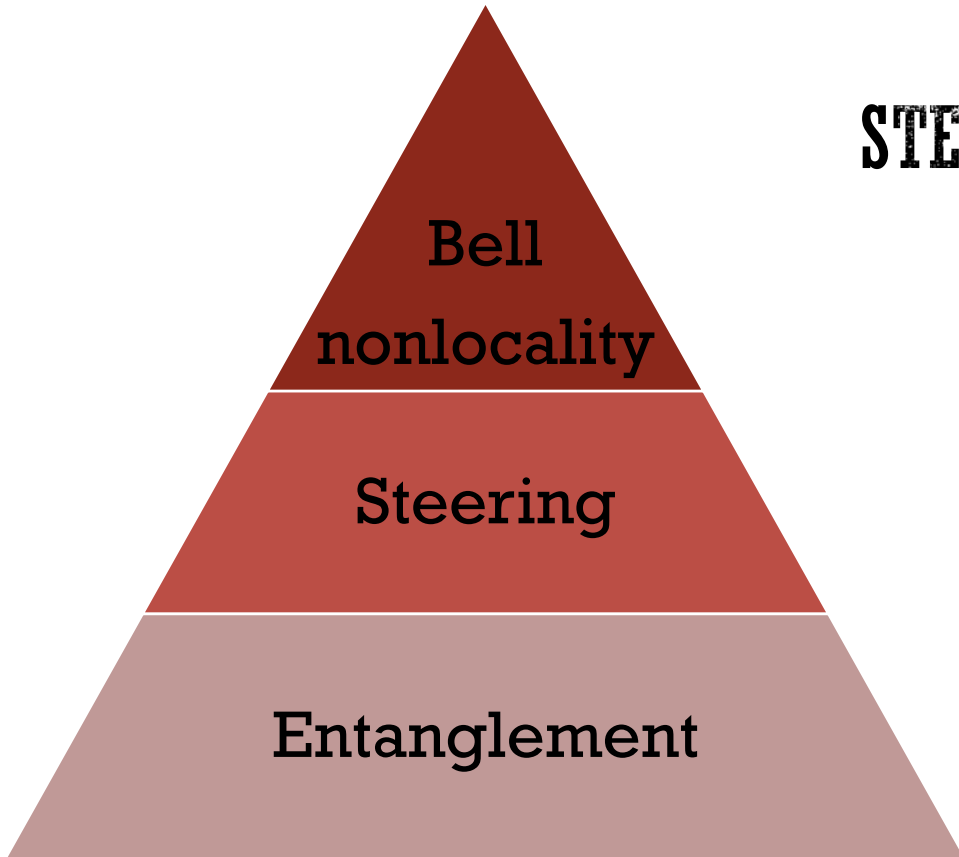
Failure of Hybrid LHV – LQS model

# ENTANGLEMENT

Failure of Local Quantum State (LQS) model

$$p(a, b|x, y) = \sum_k p_k \text{Tr}[E_{a|x} \rho_k^A] \text{Tr}[E_{b|y} \rho_k^B]$$

$$\rho = \sum_k p_k \rho_k^A \otimes \rho_k^B$$



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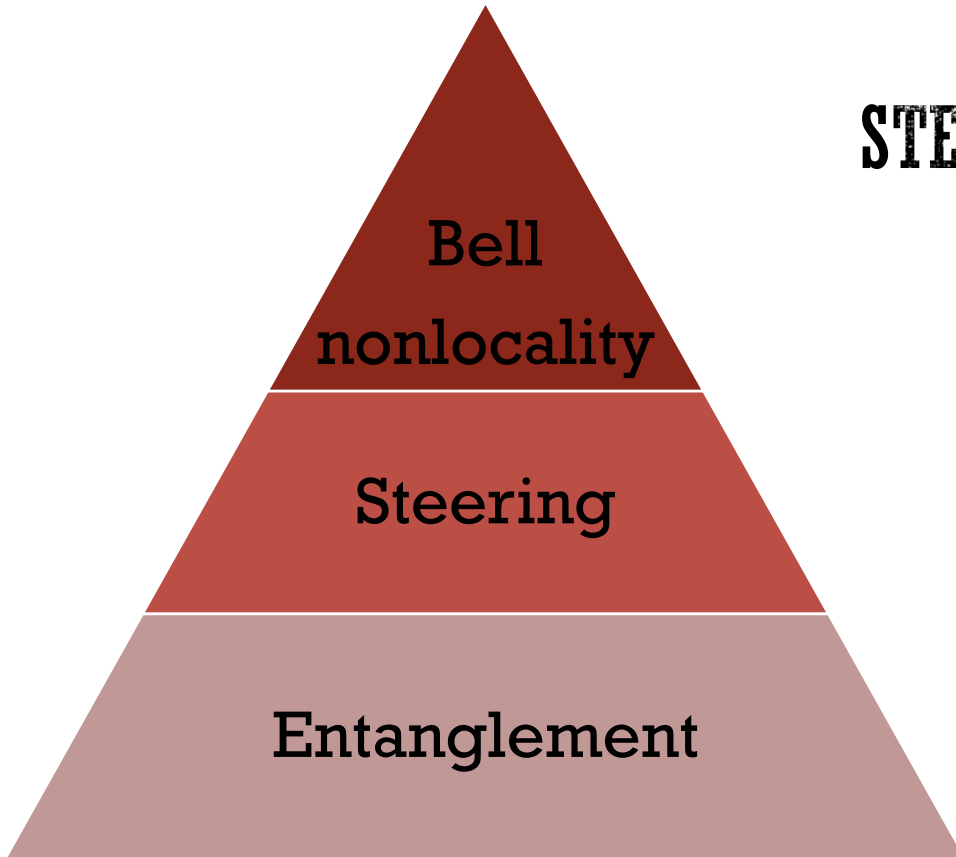
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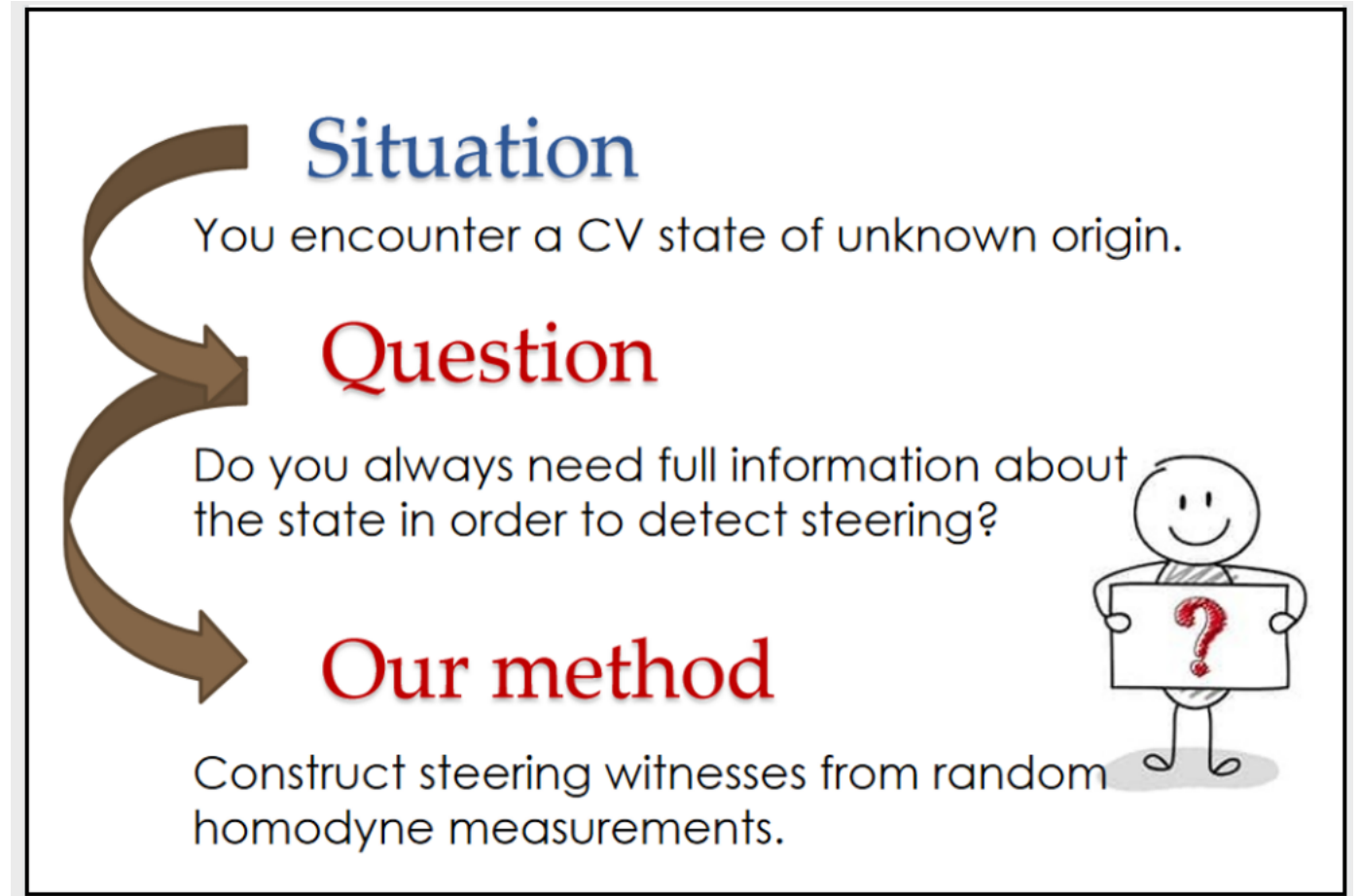
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# CONTENTS

- **Continuous Variable States**
  - ❖ Gaussian states
  - ❖ Covariance matrix (variances)
- **Gaussian Steering**
  - ❖ Witnesses method
- **Detection strategy**
  - ❑ Random measurements
  - ❑ Semidefinite programming



# CONTINUOUS VARIABLE STATES

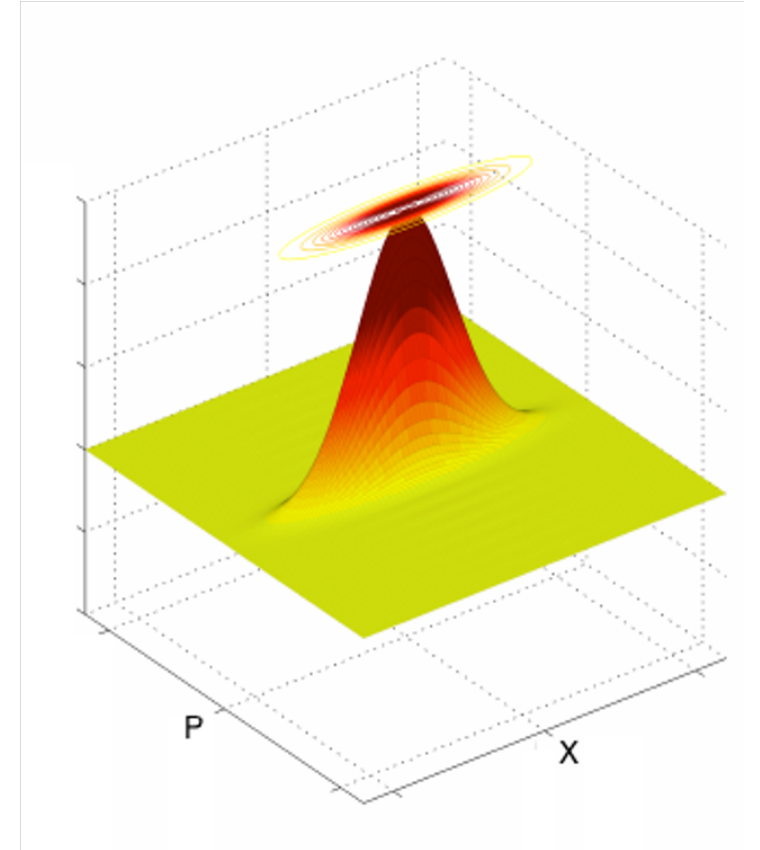
$N$  bosonic modes:  $\mathcal{H} = \bigotimes_{k=1}^N \mathcal{H}_k$

$$\hat{R}^T \equiv (\hat{R}_1, \dots, \hat{R}_{2N}) = (\hat{x}_1, \hat{p}_1, \dots, \hat{x}_N, \hat{p}_N)$$

CCM:  $[\hat{R}_i, \hat{R}_j] = i\Omega_{ij}\hat{1}, \quad i, j = 1, \dots, 2N$

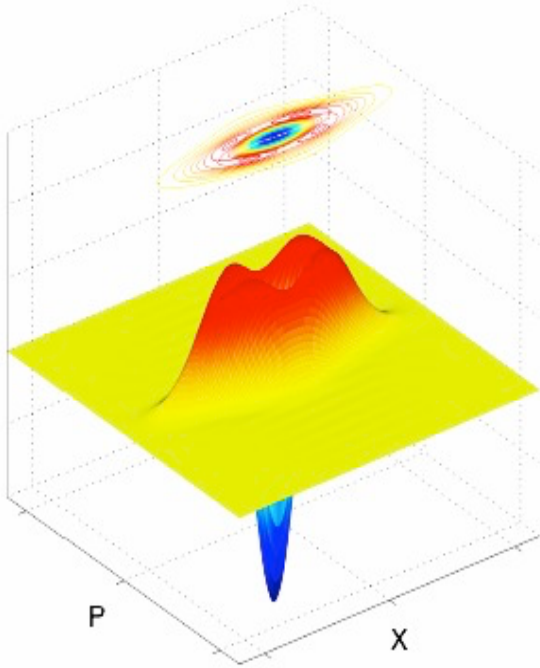
where  $\Omega_{ij} \equiv [\Omega]_{ij}$  are the elements of:

$$\Omega_N = \bigoplus_1^N \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}.$$





# FOURIER-WEYL RELATION



- Complete set of operators: Weyl displacement operators

$$\hat{D}(r) = e^{ir^T \Omega_N \hat{R}}$$

$r^T = (x_1, p_1, \dots, x_N, p_N)$  - a real vector of phase space variables.

- Fourier-Weyl relation

$$\hat{\rho} = \frac{1}{(2\pi)^{2N}} \int_{\mathbb{R}^{2N}} d^{2N} r \text{Tr}[\hat{D}^\dagger(r) \hat{\rho}] \hat{D}(r),$$

where the characteristic function is given by

$$\chi_\rho = \text{Tr}[\hat{D}^\dagger(r) \hat{\rho}]$$



# GAUSSIAN STATES

- Characteristic function with zero first moments:

$$\chi_G = e^{\frac{1}{4}r^T \Omega_N^T \gamma_{AB} \Omega_N r},$$

- Covariance matrix (CM)  $\gamma_{AB}$

$$\gamma_{ij} = \langle \{\hat{R}_i - \langle \hat{R}_i \rangle, \hat{R}_j - \langle \hat{R}_j \rangle\}_+ \rangle_\rho.$$

- Uncertainty relation:

$$\gamma_{AB} + i\Omega_N \geq 0.$$

- Bipartite CM

$$\gamma_{AB} = \begin{pmatrix} \gamma_A & \gamma_{12} \\ \gamma_{12}^T & \gamma_B \end{pmatrix}$$

- Partial Gaussian measurements

$$\text{Tr}_A[(\hat{A} \otimes \hat{I}_B)\hat{\rho}],$$

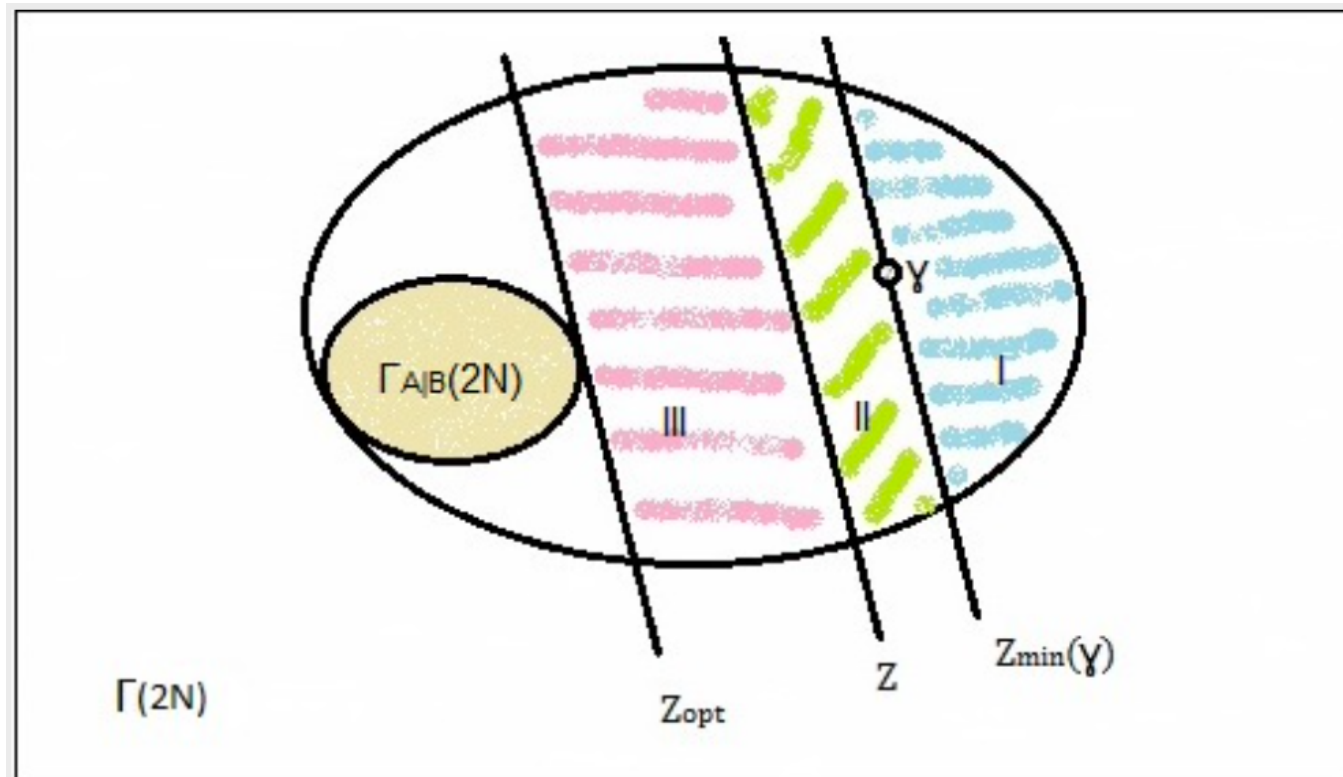
- Gaussian operator  $\hat{A}$  with CM  $T^A$
- Conditional CM after measurement

$$\gamma_B^A = \gamma_B - \gamma_{12}^T (\gamma_A + T^A)^{-1} \gamma_{12}.$$



**Theorem 2.** A bipartite quantum Gaussian state  $\rho_{AB}$  is Alice  $\rightarrow$  Bob non-steerable by means of Gaussian measurements if and only if there exists a covariance matrix corresponding to Bob's system  $\sigma_B$  satisfying  $\sigma_B + i\Omega_{N_B} \geq 0$ , such that:

$$\gamma_{AB} \geq 0_A \oplus \sigma_B.$$



## GAUSSIAN STEERING

The set of non-steerable CMs by Gaussian measurements, similarly to the set of all CMs, forms a closed convex subset of the space of all covariance matrices.

### Hahn-Banach Theorem

This allows to completely describe the set of non-steerable CMs by a family of linear inequalities representing the steering witnesses (SWs).



*Theorem 4.* A CM  $\gamma_{AB}$  of two parties consisting of  $N = N_A + N_B$  modes is Alice to Bob steerable by means of Gaussian measurements if and only if there exists a  $Z$  such that:

$$\text{Tr}[Z\gamma_{AB}] < 1,$$

where  $Z$  is a real symmetric  $2N \times 2N$  matrix satisfying

$$Z \geq 0 \quad \text{and} \quad \text{str}[Z_B] \geq \frac{1}{2},$$

where  $Z_B$  denotes the principal submatrix of  $Z$  belonging to the subsystem of Bob. Matrices  $Z$  are called steering witnesses based on second moments.

## STEERING WITNESSES

str – stands for the "symplectic trace" which is the sum of singular eigenvalues of the matrix.

We are going to construct the witness from the measurements.





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$$Z_B + i \frac{1}{2N_B} \Omega_B \geq 0$$

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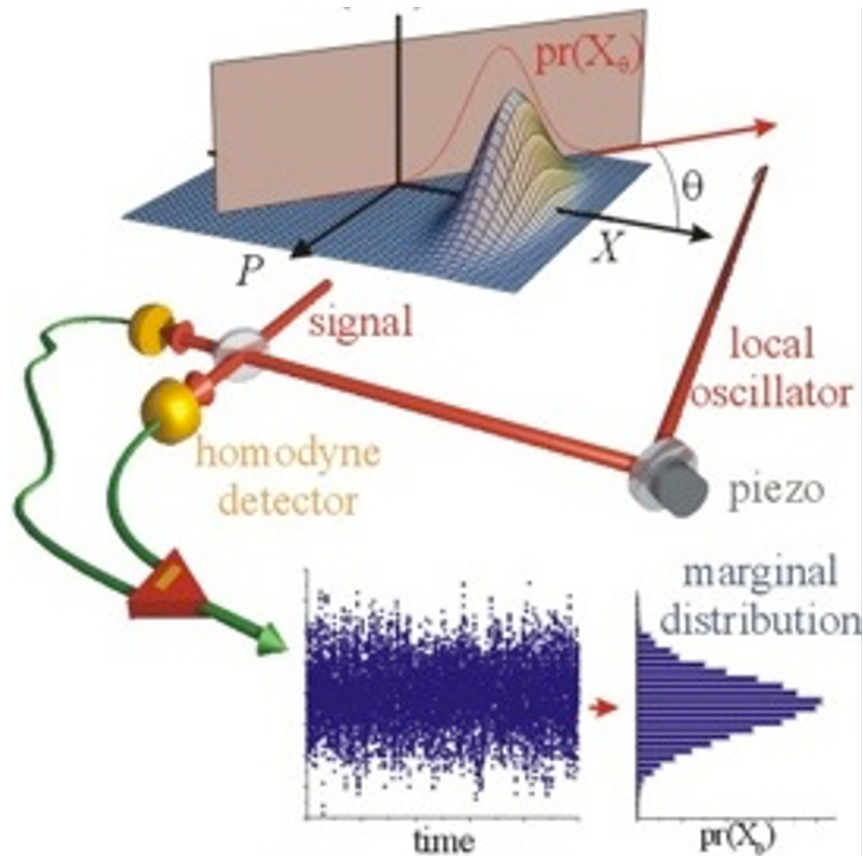
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# HOMODYNE DETECTION



- Measurement direction  $\theta$

- Quadrature to be measured

$$\hat{X}_\theta = \hat{x} \cos \theta + \hat{p} \sin \theta$$

- Marginal probability distribution  $Pr(\hat{X}_\theta)$

- The variance

$$\sigma^2 = \langle \hat{X}_\theta^2 \rangle - \langle \hat{X}_\theta \rangle^2 = Tr[P\gamma]$$

where

$$P = \begin{pmatrix} \cos^2 \theta & \cos \theta \sin \theta \\ \cos \theta \sin \theta & \sin^2 \theta \end{pmatrix}$$



# DETECTION SCHEME

- For two-mode states

$$\hat{k} = \exp(i\varphi) \cos \phi \hat{a} + \sin \phi \hat{b}$$

- Generalized quadrature

$$\hat{x}_\theta = \frac{\exp(-i\theta)\hat{k} + \exp(i\theta)\hat{k}^\dagger}{\sqrt{2}}$$

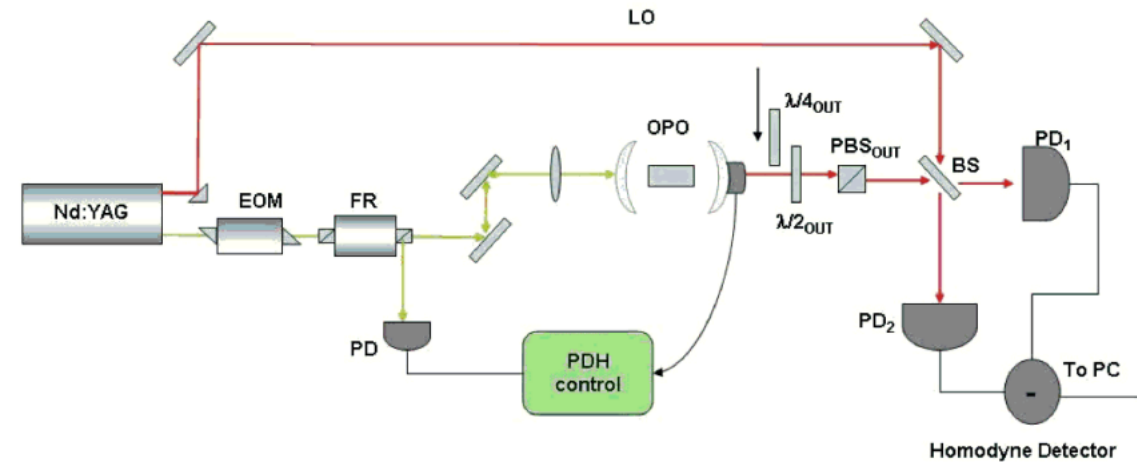


FIG. 1: Experimental setup: A type-II OPO containing a periodically poled crystal (PPKTP) is pumped by the second harmonic of a Nd:YAG laser. At the OPO output, a half-wave plate ( $\lambda/2_{\text{out}}$ ), a quarter-wave plate ( $\lambda/4_{\text{out}}$ ) and a  $\text{PBS}_{\text{out}}$  select the mode for homodyning. The resulting electronic signal is acquired via a PC module.



# OPTIMIZATION ALGORITHM

$$\begin{aligned} \text{minimize}_x : & \quad \mathbf{c} \cdot \mathbf{m} \\ \text{subject to :} & \quad Z = \sum_i c_i P_i \\ & \quad Z \geq 0 \\ & \quad Z_A + ix\Omega_1 \geq 0 \\ & \quad Z_B + i\left(\frac{1}{2} - x\right)\Omega_1 \geq 0 \end{aligned}$$

where  $\mathbf{m} = \text{Tr}(\mathbf{P}\gamma)$ , with  $\mathbf{P}$  the vector of measurements  $P_i$ .





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Performe a random  
measurement



# OPTIMIZATION ALGORITHM

$$\begin{aligned} \text{minimize}_x : & \quad \mathbf{c} \cdot \mathbf{m} \\ \text{subject to :} & \quad Z = \sum_i c_i P_i \\ & \quad Z \geq 0 \\ & \quad Z_A + ix\Omega_1 \geq 0 \\ & \quad Z_B + i\left(\frac{1}{2} - x\right)\Omega_1 \geq 0 \end{aligned}$$

where  $\mathbf{m} = \text{Tr}(\mathbf{P}\gamma)$ , with  $\mathbf{P}$  the vector of measurements  $P_i$ .

Performe a random  
measurement

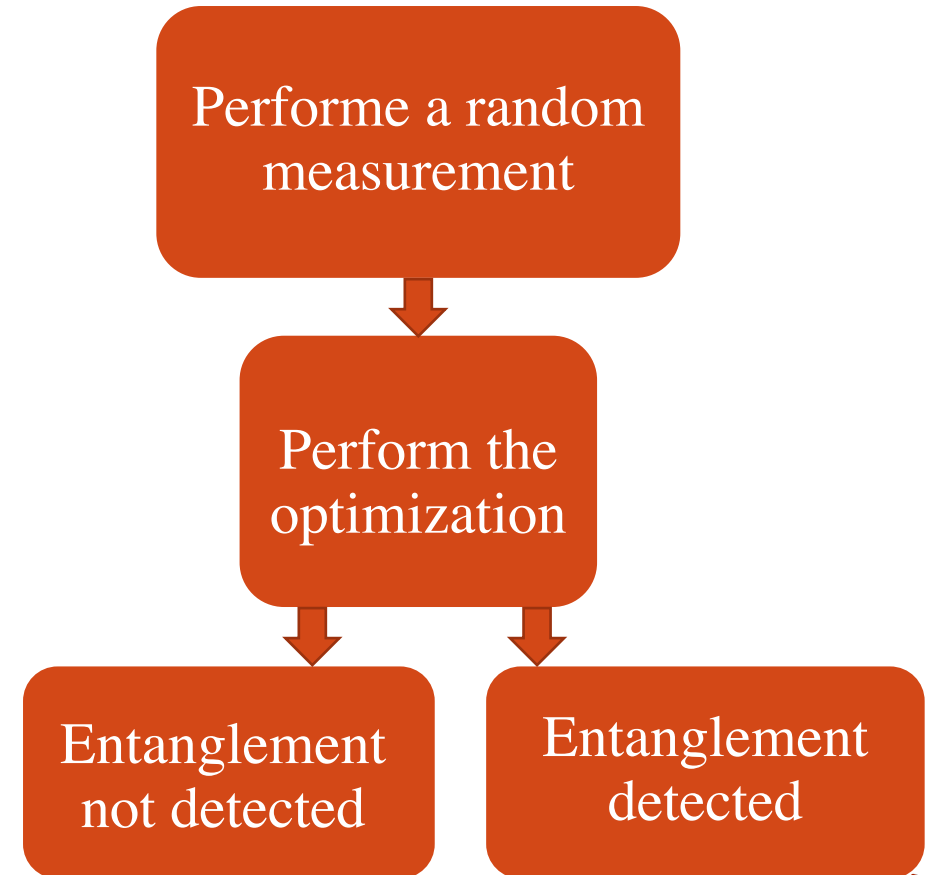
Perform the  
optimization



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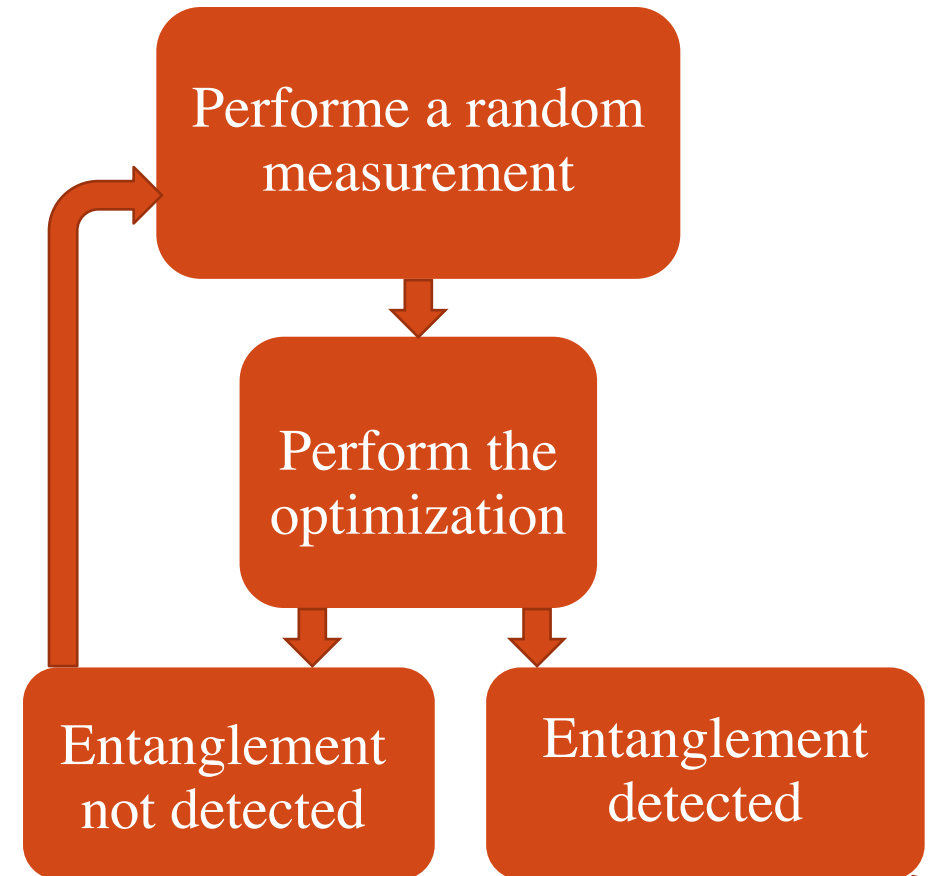
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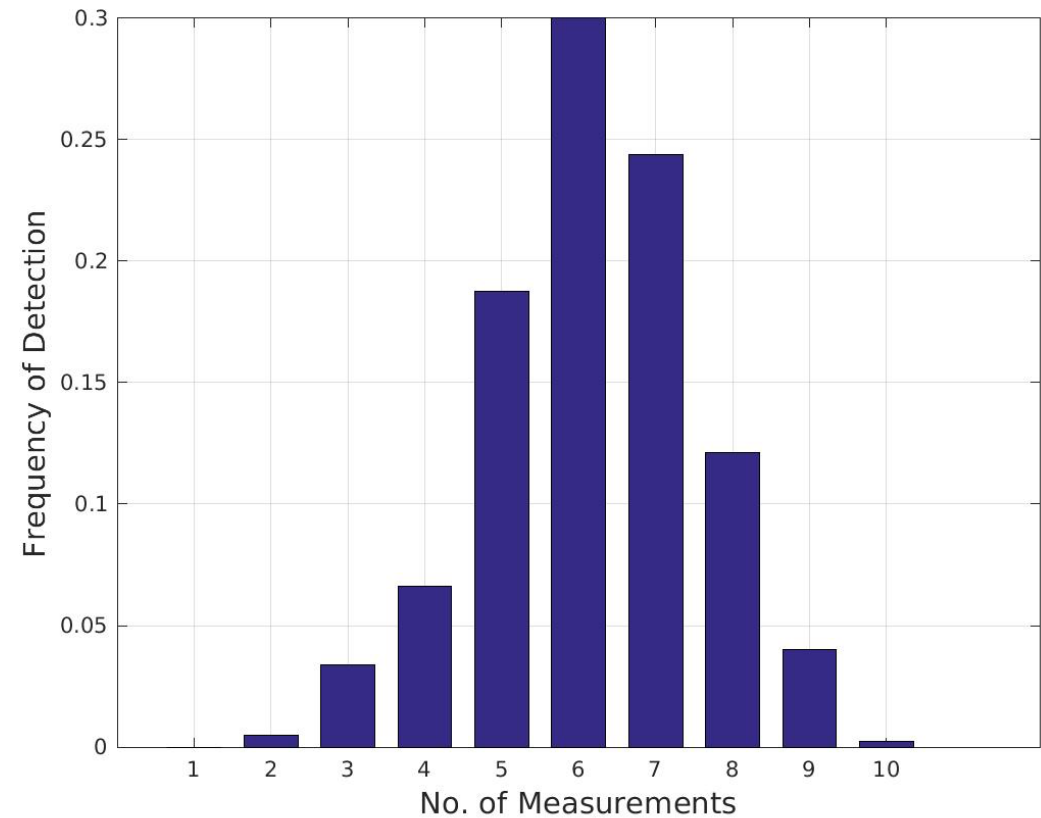
where  $\mathbf{m} = \text{Tr}(\mathbf{P}\gamma)$ , with  $\mathbf{P}$  the vector of measurements  $P_i$ .

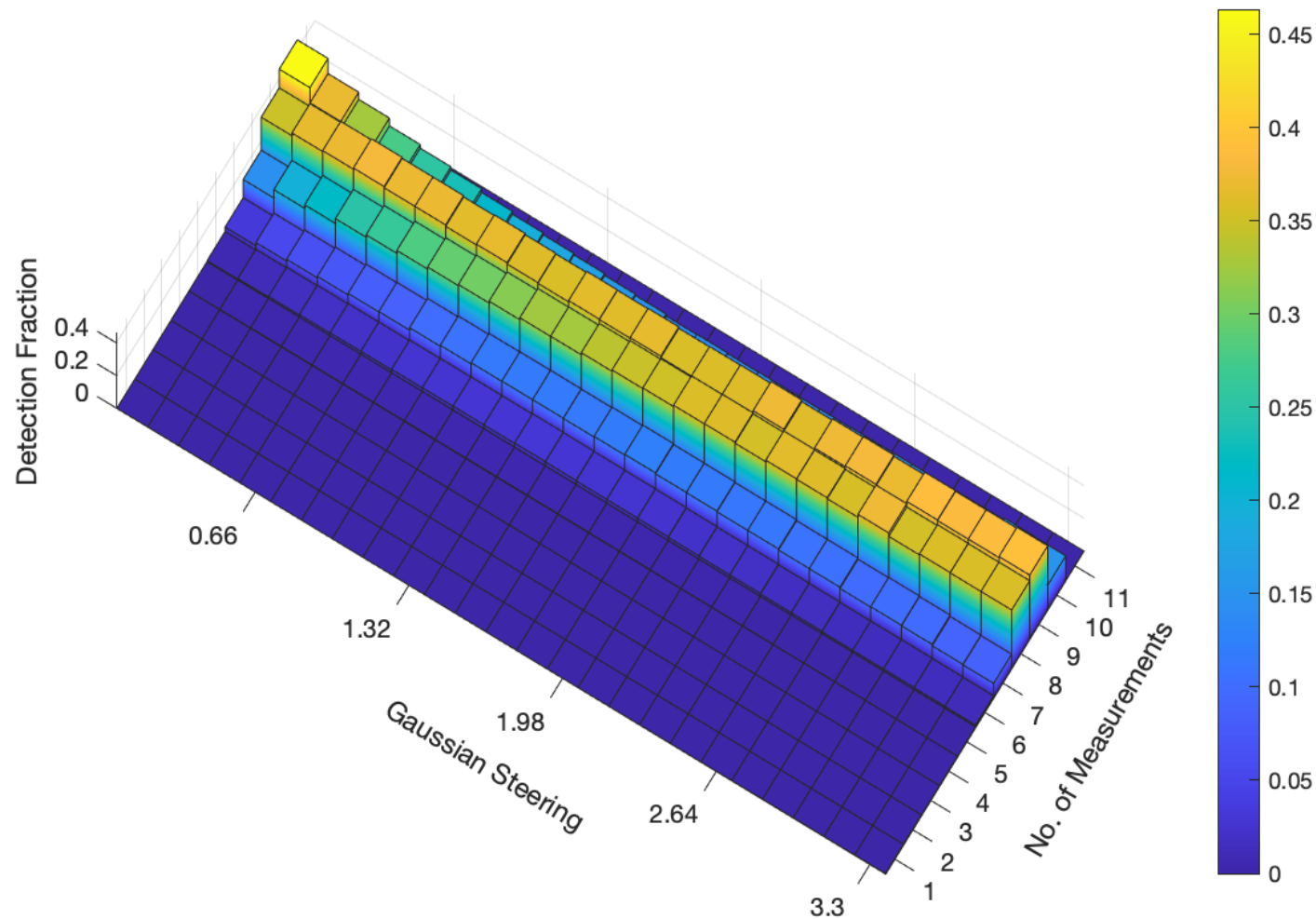




# FIXED VALUE OF STEERING

- Two-mode covariance matrix
- Full tomography: 10 independent Measurements
- The data is normalized such that it sums up to 1.





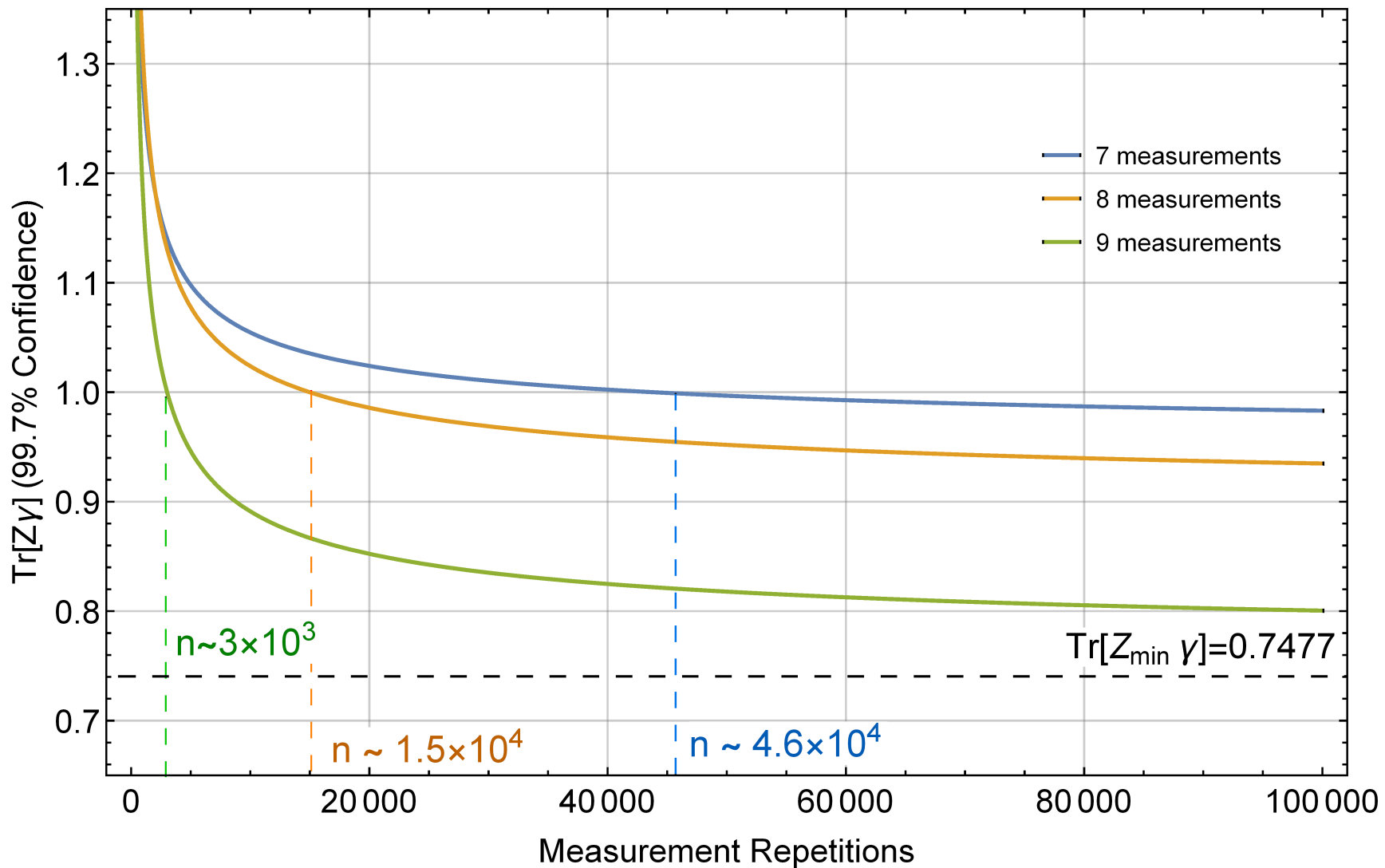
# DETECTION IN RANDOM COVARIANCE MATRICES

Fraction of steering detection for two-mode randomly generated CMs.

It represents  $5 \times 10^5$  runs of the algorithm where measurement directions are added successively and SW is evaluated at every round until steering is certified.

The data is normalised for every value of steering such that they sum up to 1.





$$\Delta \bar{Z} = \sqrt{\frac{2}{n-1}} \sqrt{\sum_i c_i^2 m_i^2}.$$

## STATISTICAL ANALYSIS

The statistical estimate of  $Z^-$  with the maximum of  $3\sigma$  confidence interval.

The horizontal black dashed line indicates the minimal value of the witness for the considered CM  $\text{Tr}[Z_{\min} Y] = 0.7477$ .

The vertical dashed lines indicate the number of measurement repetitions required to detect Gaussian steering with 7 (blue), 8 (orange) and 9 (green) measurement settings.





**THANK YOU FOR  
ATTENTION!**

